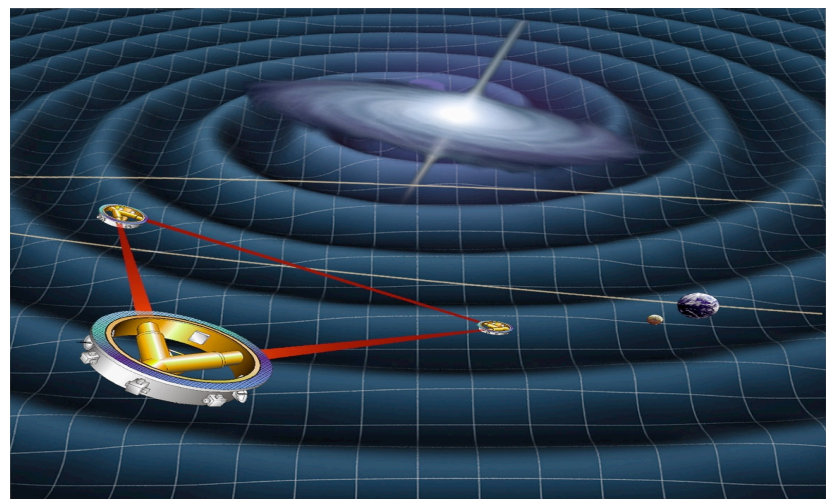


Black Holes

Kip Thorne



Chandrasekhar Centennial Symposium
University of Chicago, 15 September 2010

What I Will NOT Cover

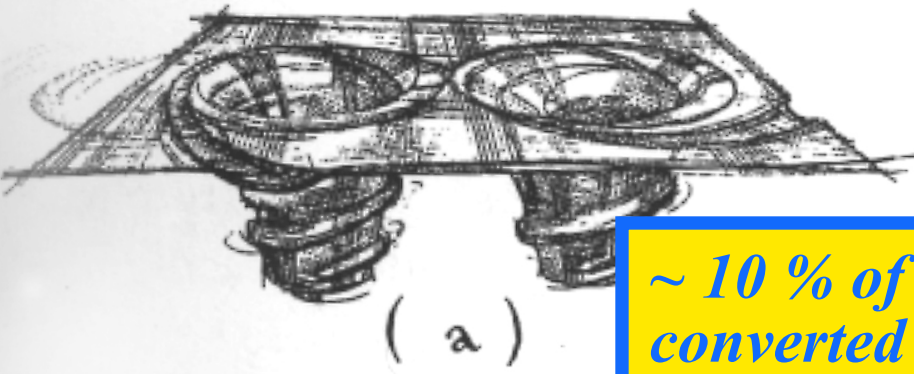
- Semi-Classical Aspects of Black Holes
- Quantum Aspects of Black Holes
- Higher-Dimensional Black Holes
- Astrophysics of Black Holes

[talk by Priyamvada Natarajan tomorrow]

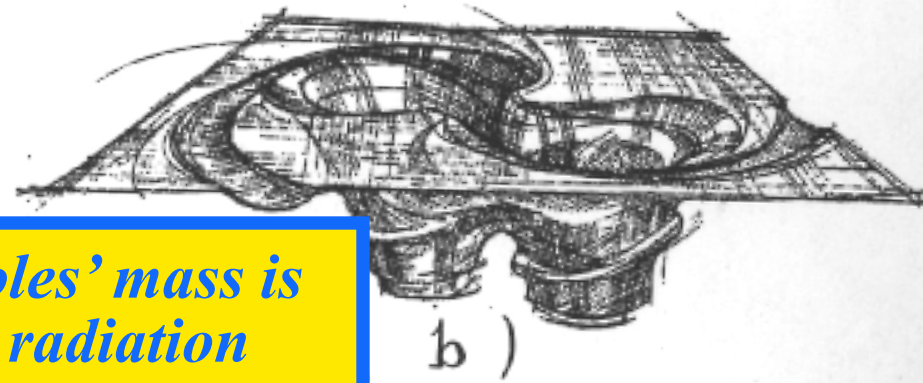
What I Will Cover

- Black Holes in the context of
 - » Numerical Relativity
 - » Gravitational Wave Observations
- A new “golden age” of black-hole research

The “Holy Grail”: Collisions of Black Holes - The most violent events in the Universe



(a)



b)

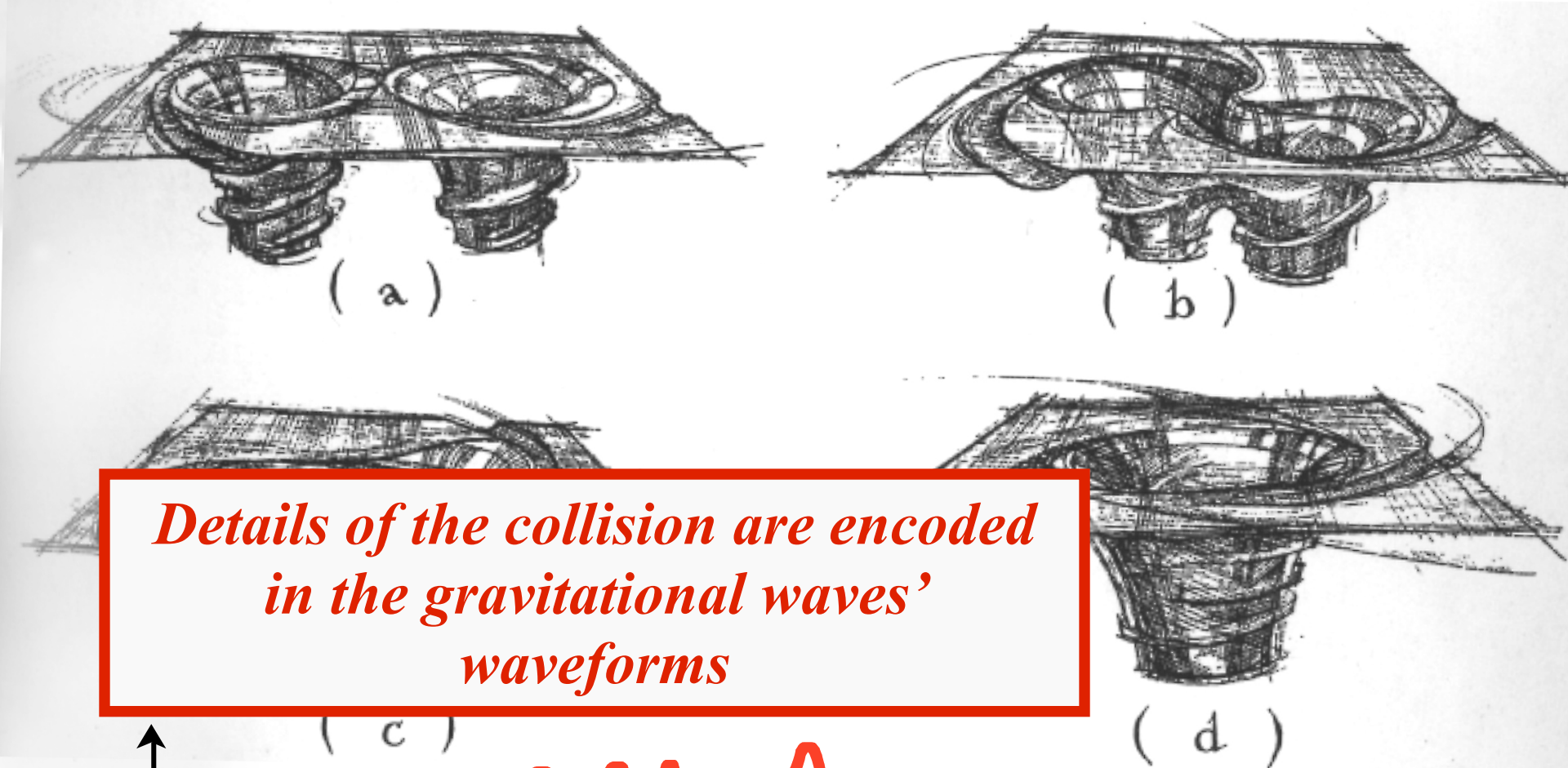
*~ 10 % of holes' mass is converted to radiation
[contrast with nuclear fusion: < 0.5 %]*



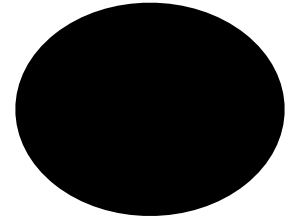
*GW Luminosity $\sim 0.1 Mc^2 / (100 GM/c^3)$
 $= 0.001 c^2/G \sim 10^{24} L_{\text{sun}} \sim 10^4 L_{\text{EM universe}}$*

*No Electromagnetic Waves emitted whatsoever
- except from, e.g. disturbed accretion discs*

Collisions of Black Holes: The most violent events in the Universe



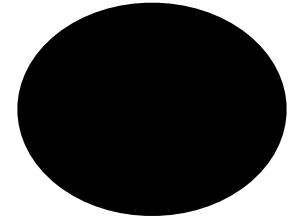
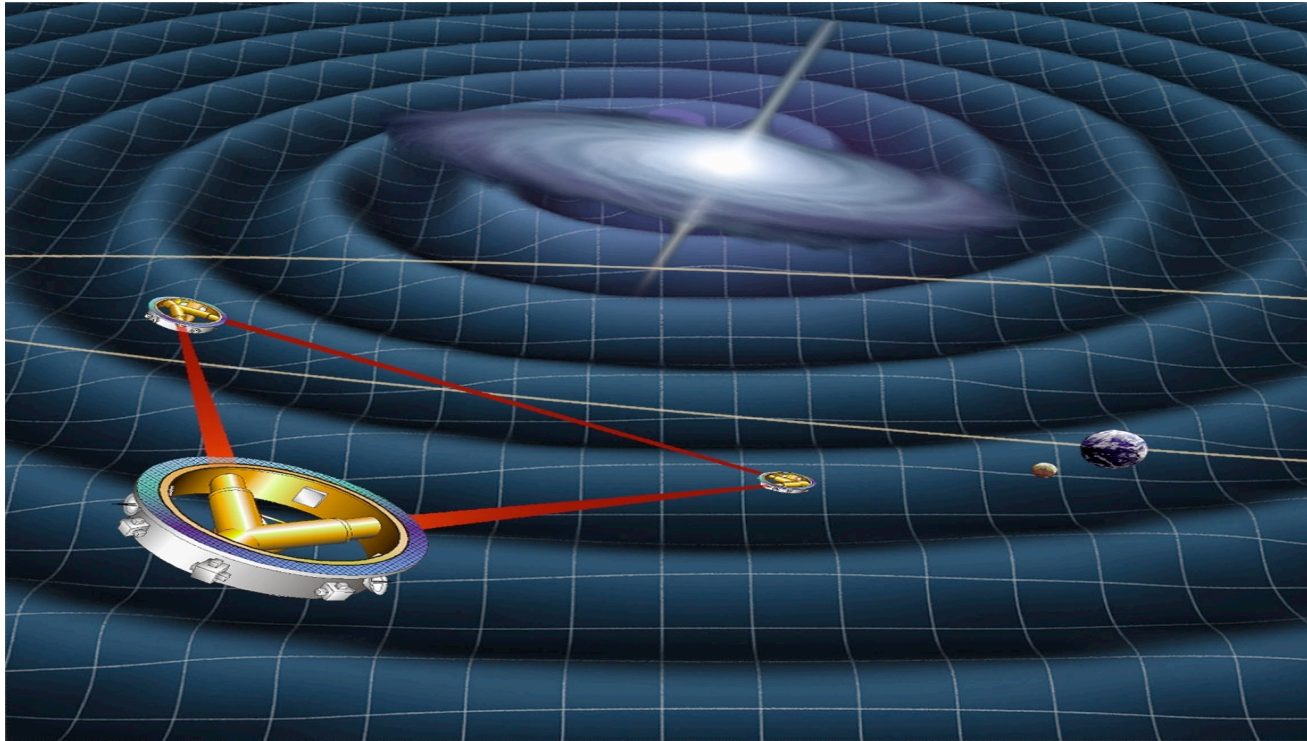
The Context: *LIGO*



Black Hole Masses:
~2 to 1000 M_{sun}

- ~700 scientists, ~ 60 Institutions, 12 nations,
- Initial detectors and searches: 2005-2010
- Advanced detectors: Begin installation next Wednesday
- All data analyzed jointly with VIRGO

The Context: *LISA*

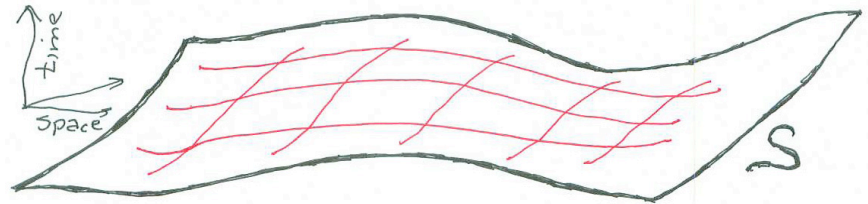


Black Hole Masses:
~10,000 M_{sun} to
~10 million M_{sun}

- Joint NASA / ESA Mission
- LISA Pathfinder (technology test) mission: 2013
- LISA launch: ~2022?

Numerical Relativity: How is it Done?

- Evolve the geometry of spacetime - not fields in spacetime
- Choose an initial spacelike 3-dimensional surface S
 - » Put a coordinates on S



- Specify: 3 -metric g_{ij} and *Extrinsic Curvature* K_{ij} of S
 - » Subject to *constraint equations* [analogues of $\text{Div } B = 0$]
- Lay out coordinates to future by specifying Lapse function α and Shift function β^i
- Integrate 3 -metric forward in time via *dynamical equations*
$$ds^2 = -\alpha^2 dt^2 + g_{ij} (dx^i - \beta^i dt) (dx^j - \beta^j dt)$$

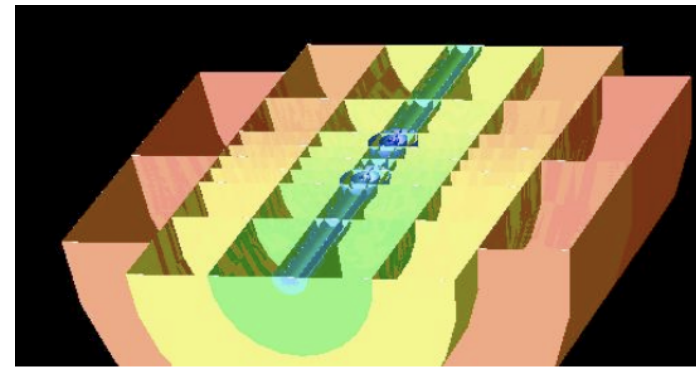
Numerical Relativity

Two Major Pitfalls

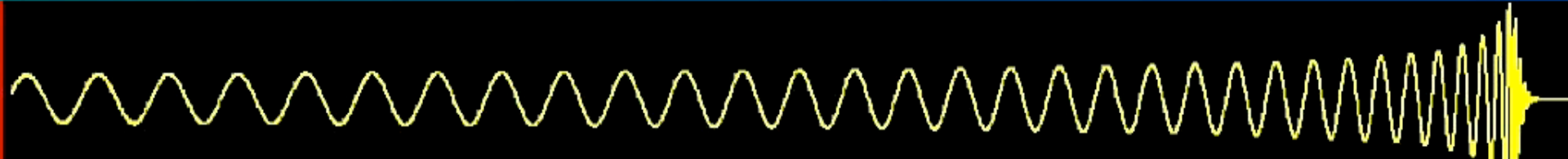
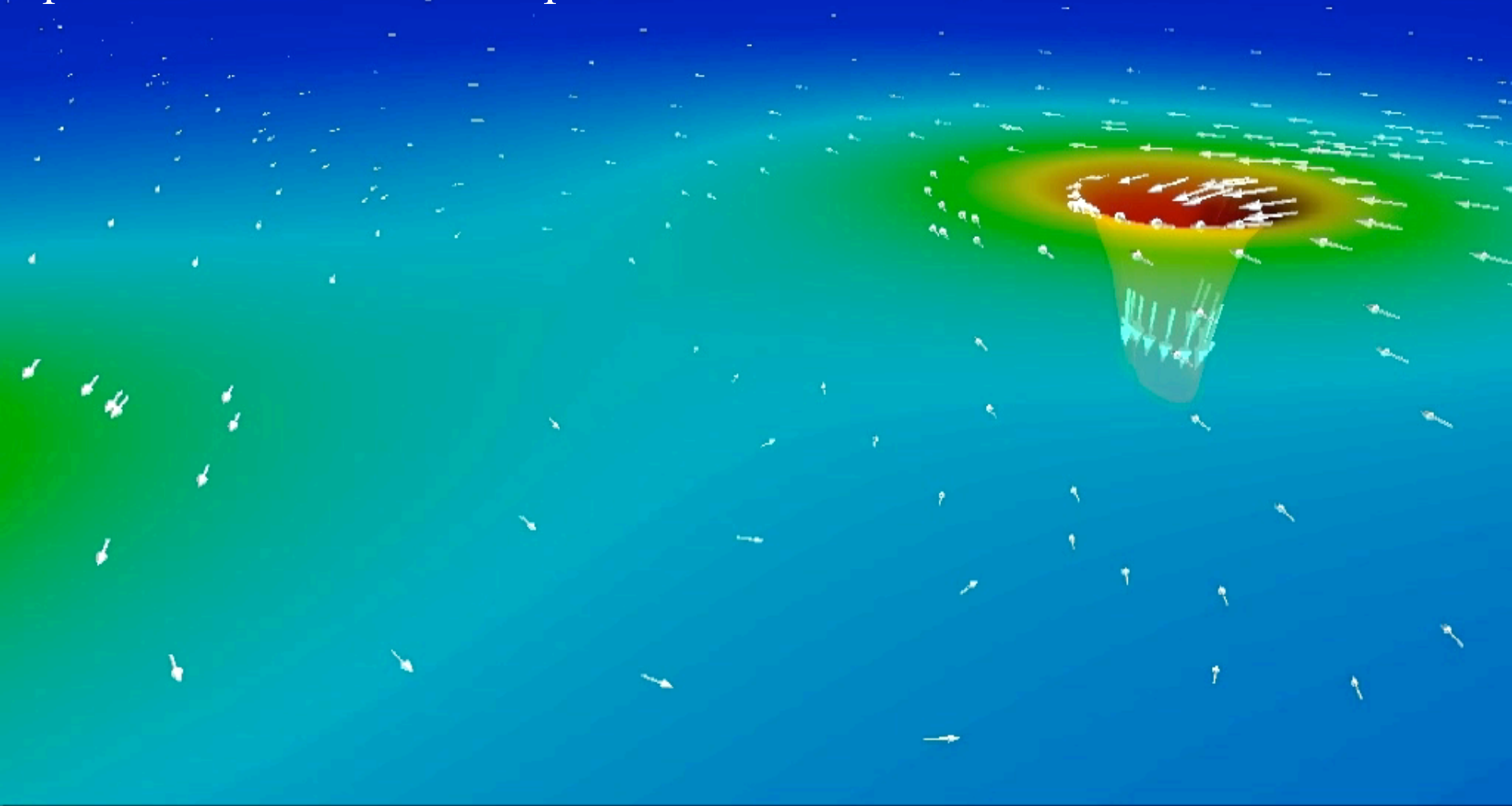
- Constraint-Violation Instabilities:
 - » Slight initial error in constraints (analog of $\text{Div } B = 0$) blows up in time
 - » Solved after ~ 5 years of struggle
- Coordinates become singular [“gauge instability”]
 - » Only recently solved *robustly* in highest-precision codes

Two Mature Approaches

- Finite-difference
 - » Robust, power-law convergence
- Spectral
 - » More complicated, less robust
- but exponential convergence



Caltech/Cornell/CITA - Kidder, Pfeiffer, Scheel,
Teukolsky, Lindblom, ...
Spectral Einstein Code: SpEC



Current NR BBH Capabilities

- Many groups can simulate generic black holes (unequal masses, spins with random orientations)
 - » Princeton (Pretorius), Rochester Institute of Technology (Campanelli, ...), Goddard Spaceflight Center (Centrella, ...), U. Illinois (Shapiro, ...), Albert Einstein Institute & LSU (Pollney, ...), U. Jena (Bruegmann, ...), Georgia Tech (Laguna, ...), U. Texas (Matzner, ...), U. Maryland (Tiglio, ...), Florida Atlantic U. Tichy, ...), Barcelona (Sperhake, ...), Cornell/Caltech/CITA (Teukolsky, ...), ...

Current NR BBH Capabilities

- Mass ratios:
 - » most groups: up to 6:1
 - » Highest [RIT - Campanelli, Lousto, Zlochower]: 100:1
- Spin magnitudes, $\chi=S/M^2$
 - » most groups: up to 0.6
 - » Highest for binary: 0.95 [Cornell/Caltech/CITA, Lovelace, Scheel, Szilagyi]
 - standard method for solving constraint equations [Bowen/York] limited to $\chi < 0.93$
 - new method: superposed Kerr-Schild

“S-Matrix” Insights from NR

- S-Matrix:

- » Initial state: Masses, Spins, Orbit

- » Final state: Mass, Spin, Kick Velocity,

Gravitational Waveforms

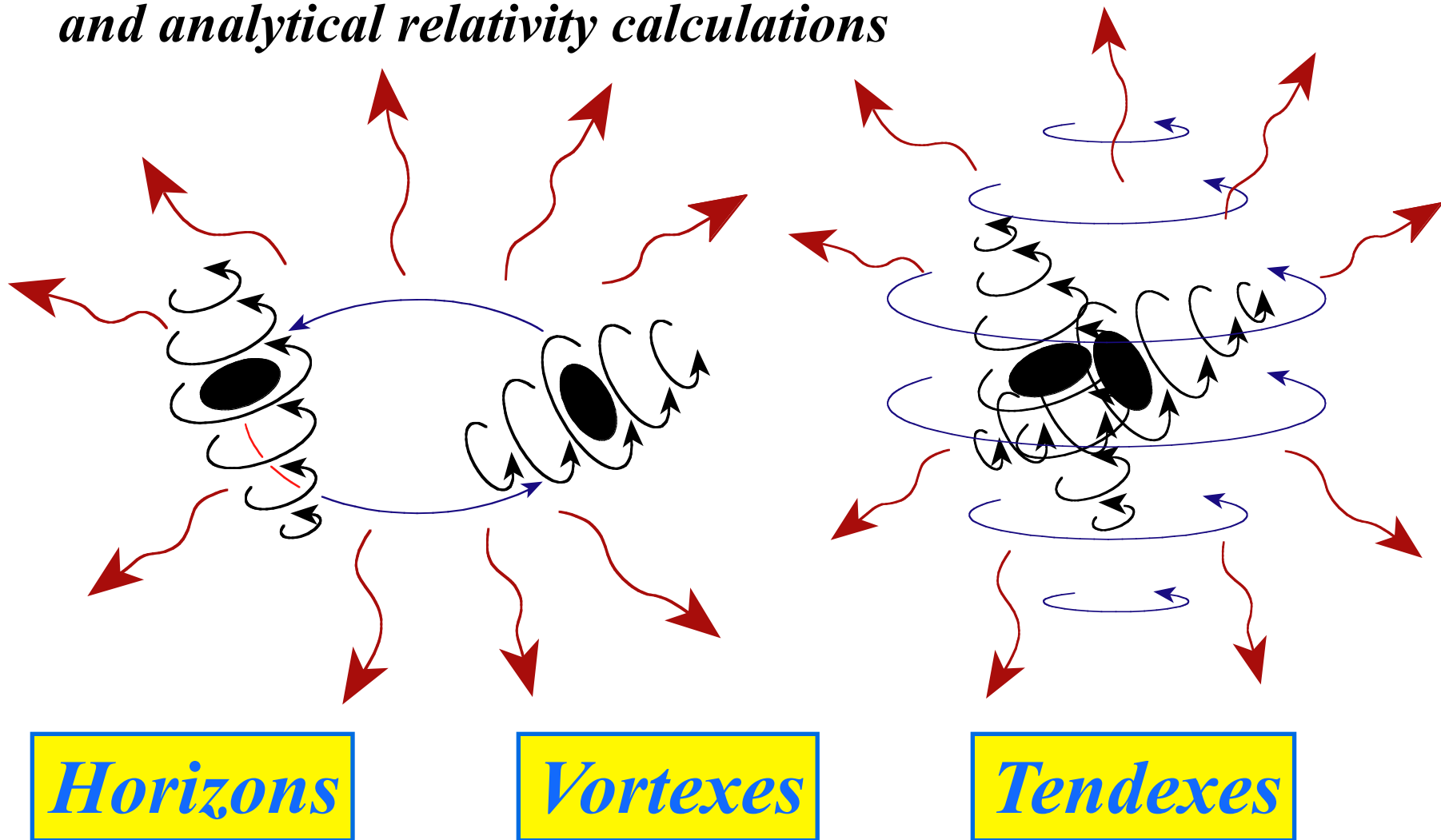
BBH Waveforms for LIGO

Clifford Will's Talk

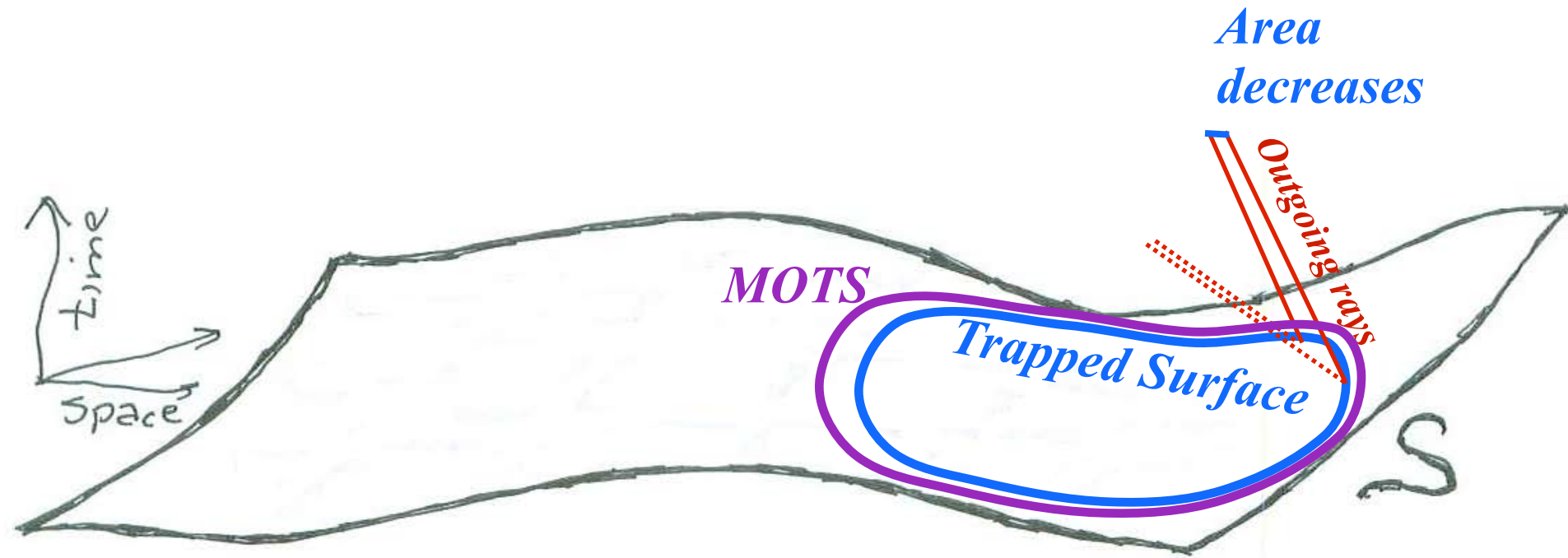
- Waveforms needed in analytic form [quickly computable, “on the fly”] for searches via “matched filter method”
 - » Accuracy needed [Lindblom, Owen, brown]: $\delta\phi \sim \delta A/A \sim 0.1$ for searches, ~ 0.01 for information extraction
 - » PN at early times; NR at late times; match
 - » NR simulations
 - 7 parameters: $M_2/M_1, S_2, S_1$
 - need ~ 1000 simulations
 - » Use PN and NR to “tune” parametrized analytical formulae, e.g. “Effective One Body”
- Underway: Large collaboration led by Buonanno [U Md] - most of the world's NR groups + several Analytical Relativity groups

Nonlinear Dynamics of Curved Spacetime

What we are learning from combined NR simulations and analytical relativity calculations



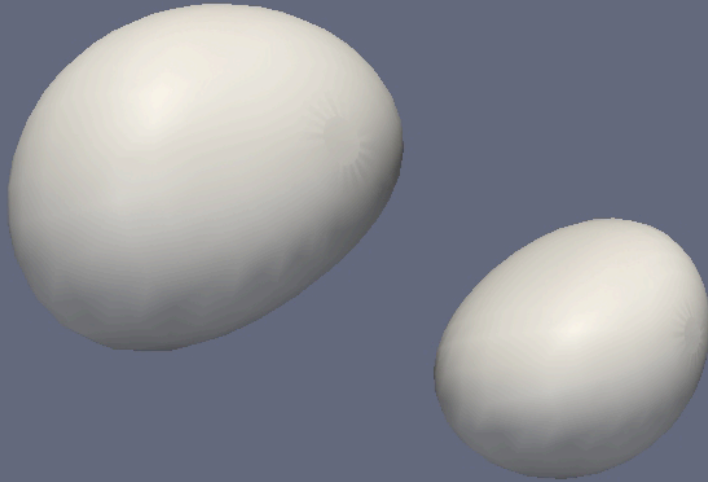
Marginally Outer Trapped Surface (MOTS) and Apparent Horizon [Penrose]



Outermost MOTS: Apparent Horizon

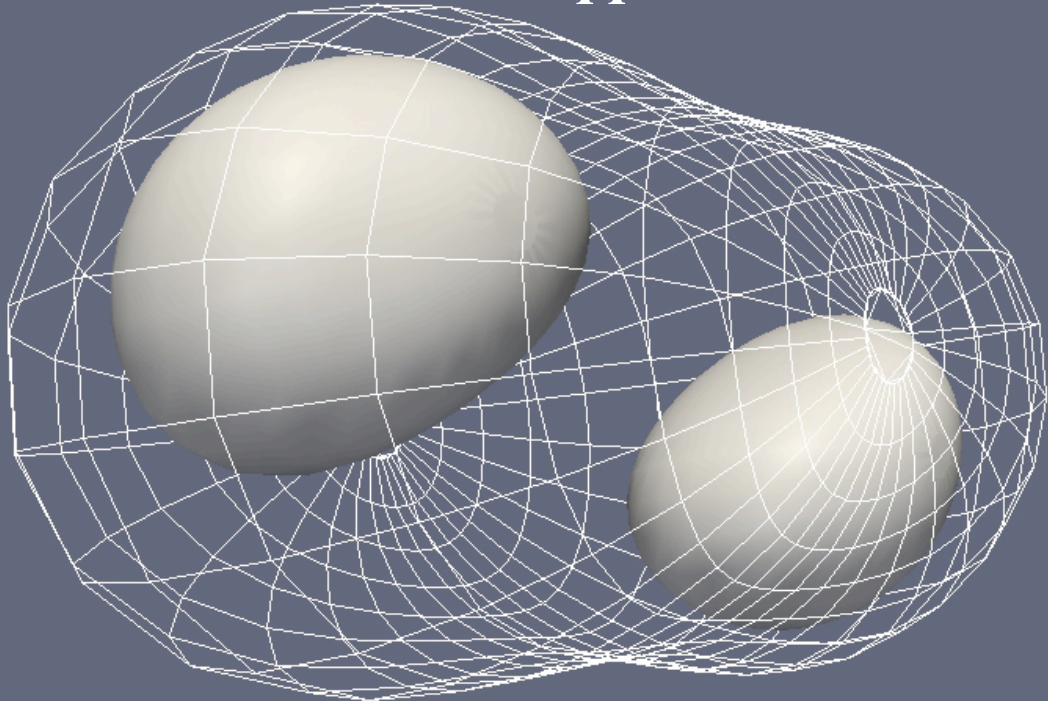
From SpEC Simulation

Apparent Horizons



Rob Owen [Cornell]

Common Apparent Horizon

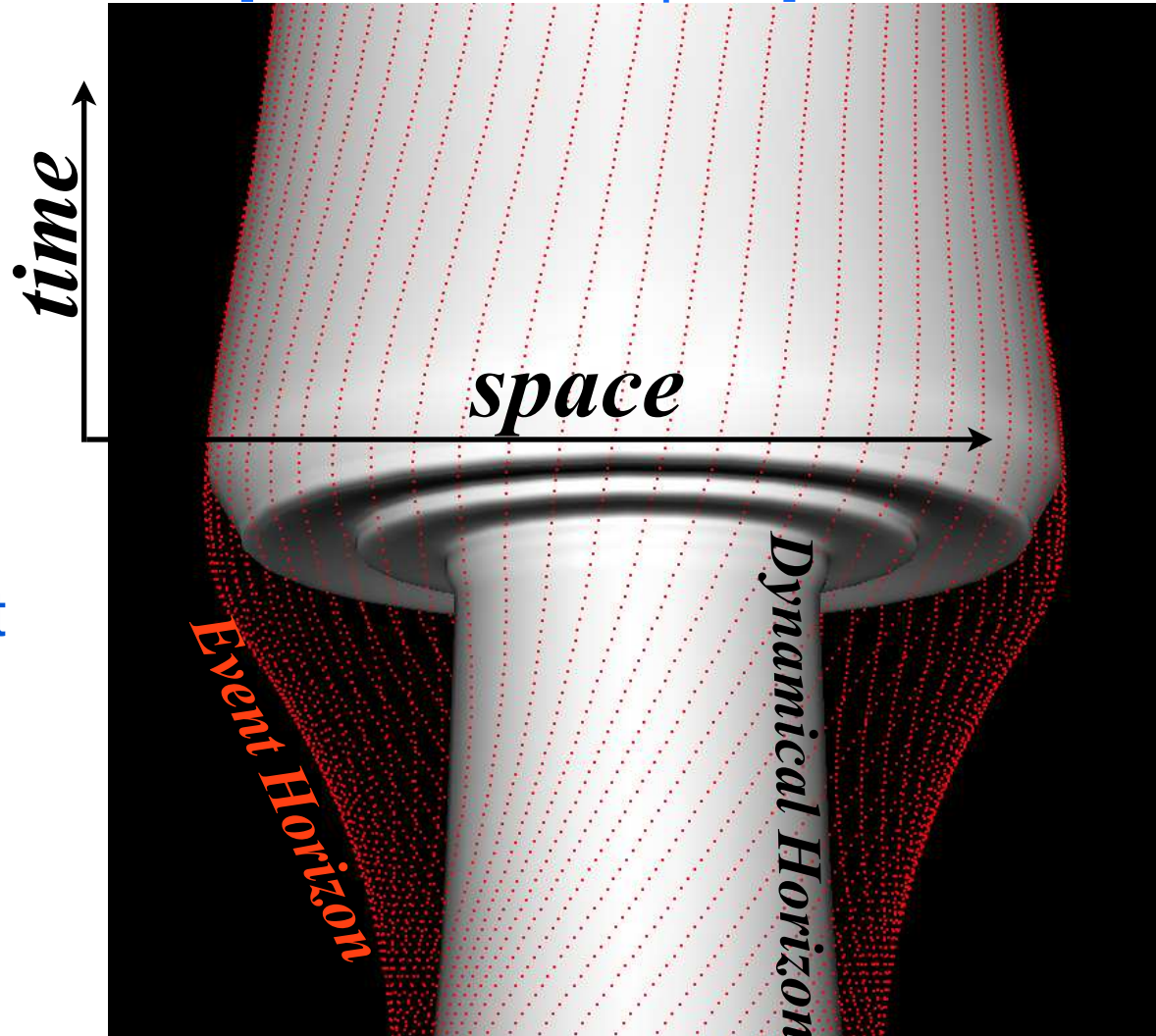


Nested MOTSs

Apparent Horizon vs Event Horizon

- Tony Chu, Harald Pfeiffer, Michael Cohen:
 - » Kerr black hole perturbed by a very strong pulse of ingoing gravitational waves [evolved with SpEC]

- *Dynamical horizon* [Ashtekar, Krishnan]: spacelike world tube of marginally outer trapped surfaces (MOTS)
- *Apparent horizon* [Penrose]: outermost MOTS
- *Event Horizon* [Hawking]: boundary of communication



Spinning Black Holes

Rochester Institute of Technology:
Campanelli, Lousto, Zlochower



Simulation:

Manuela Campanlli

Carlos Lousto

Yosef Zlochower

Visualization:

Hans-Peter Bischof

CCRG

RIT

Copyright - CCRG - 2009



Spinning Black Holes

Rochester Institute of Technology:
Campanelli, Lousto, Zlochower



Weak gravity:

$$ds^2 = -(1 - 2M/r)dt^2 + dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta (d\phi - \omega dt)^2$$

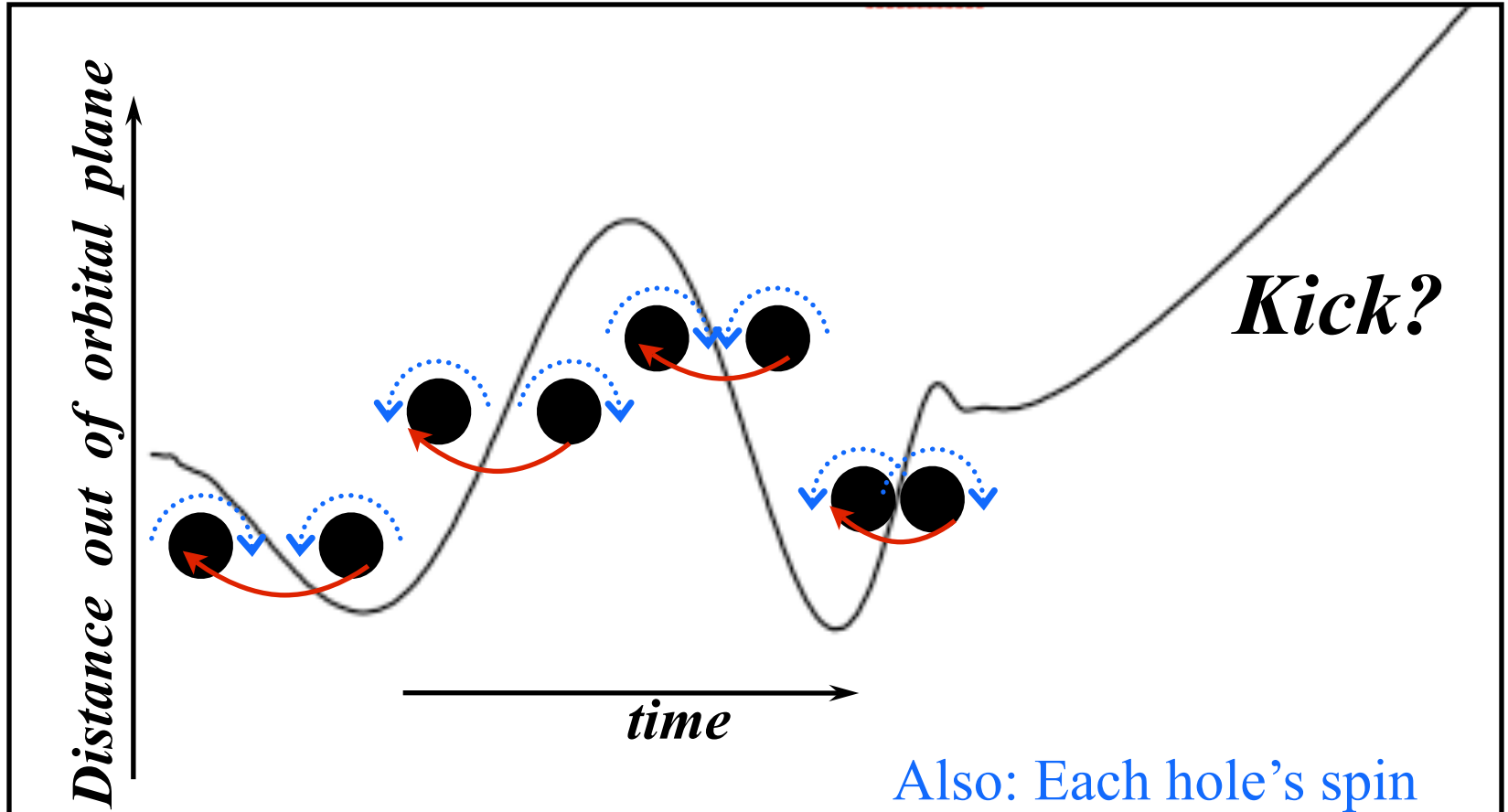
$$\omega = 2J/r^3$$
$$\mathbf{v}_{\text{FD}} = \omega r \mathbf{e}_{\hat{\phi}}$$

Analogous to 2 Vortices in a Fluid



Explanation of Bobbing

[Pretorius]

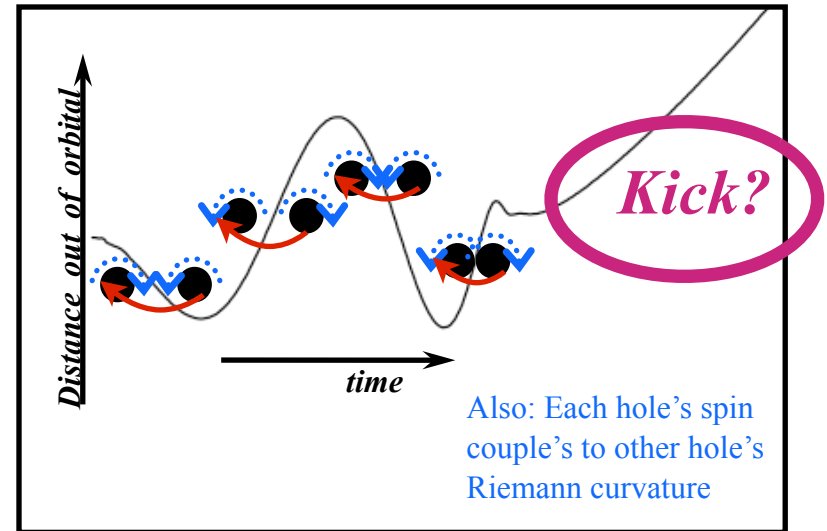


*Other new tools:
Momentum flow - Chen, Keppel, ...;
Gralla, Harte, Wald*

Also: Each hole's spin
couple's to other hole's
Riemann curvature

New Insights Into Nonlinear Dynamics of Curved Spacetime

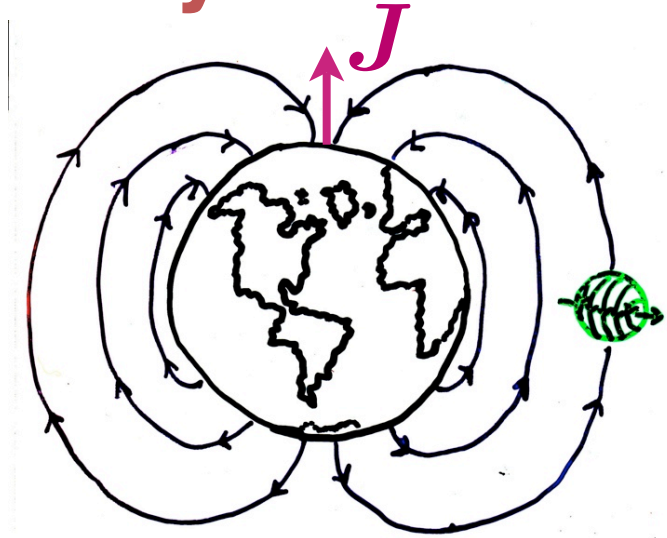
- **Cornell/Caltech/CITA Numerical Relativity Group**
 - » *Teukolsky, Kidder, Lindblom, Pfeiffer, Szilagyi, Scheel, Owen, Lovelace, Kaplan, Matthews, ...*
- **Caltech Analytical Relativity Theory (CaRT) Group**
 - » *Chen, Nichols, Thorne, ...*



Frame Dragging When Gravity is Weak

- Gyroscope precession relative to distant stars [inertial frames at “infinity”]

$$\Omega_{\text{fd}} = \frac{1}{2} \nabla \times \mathbf{v}_{\text{fd}} = - \left[\frac{\mathbf{J} - 3(\mathbf{J} \cdot \mathbf{n})\mathbf{n}}{r^3} \right]$$

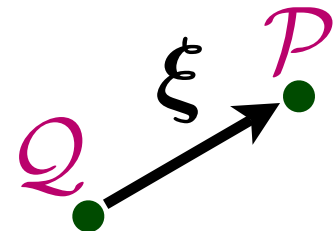


- When gravity is strong and dynamical: cannot use inertial frames at infinity.

- Differential frame dragging: Precession angular velocity at \mathcal{P} relative to inertial frames at \mathcal{Q}

$$\Delta \Omega_{\text{fd}} = \xi \cdot \nabla \Omega_{\text{fd}}$$

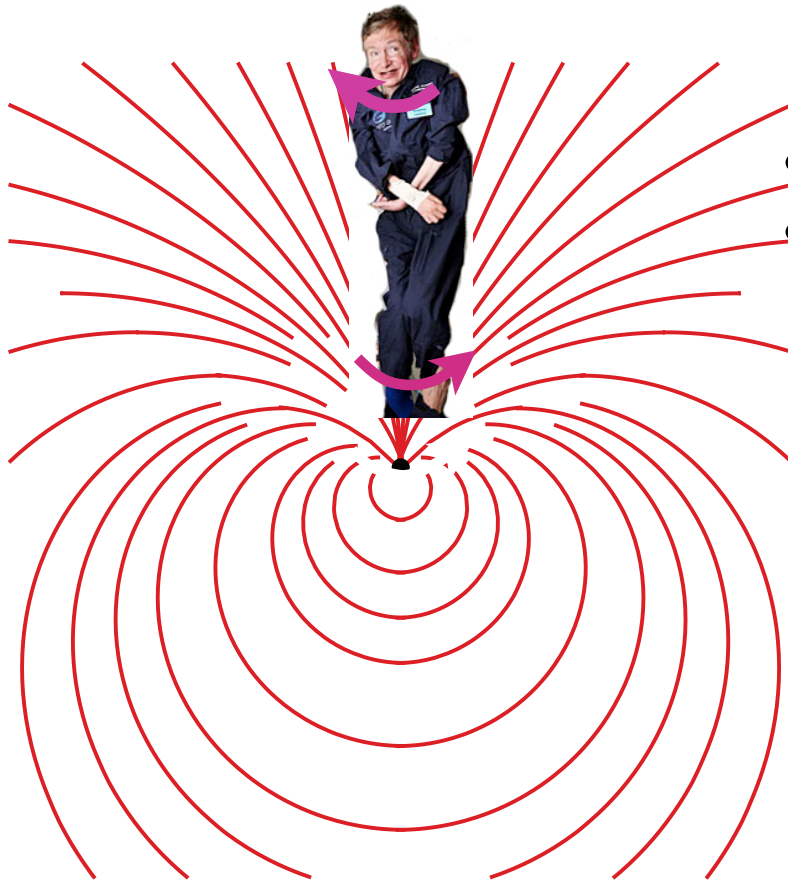
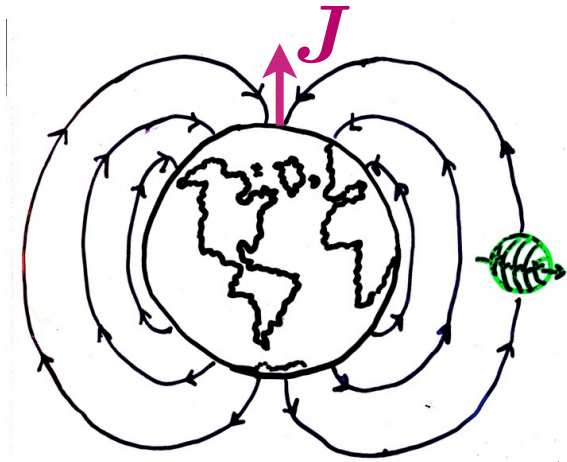
\mathcal{B} “Magnetic” part of Weyl Curvature Tensor



- \mathcal{B}_{jk} symmetric and trace free. Determined by 3 orthogonal eigendirections (principal axes) and their eigenvalues.

Frame Dragging When Gravity is Weak

- **Frame-Dragging Vortex Line:** Streamline of eigenvector n of B_{jk}
- **Vorticity of Line:** Eigenvalue B_{nn}

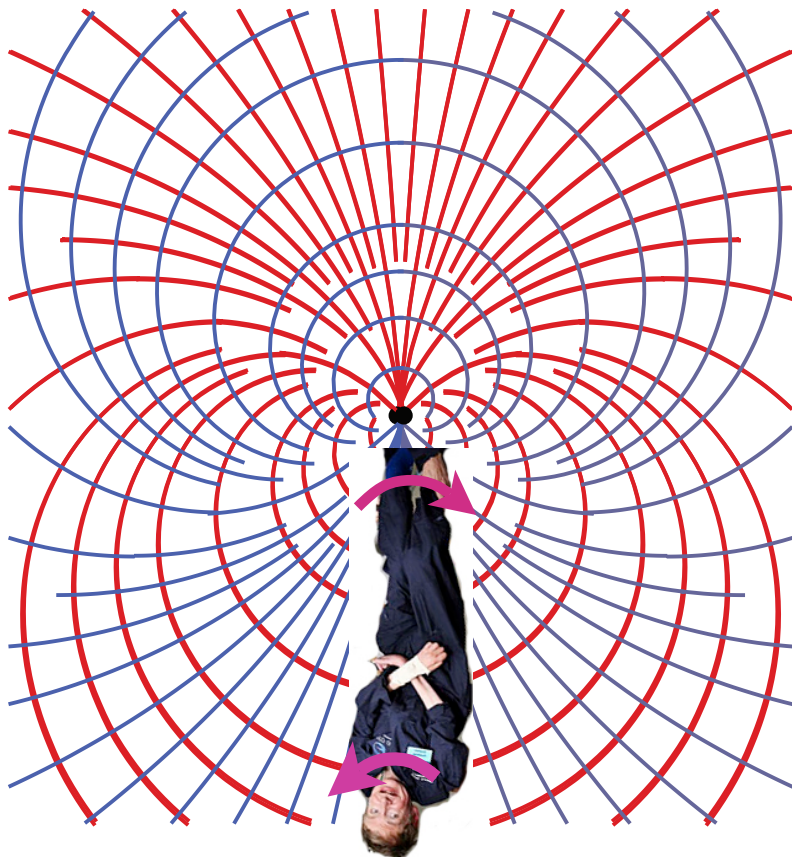
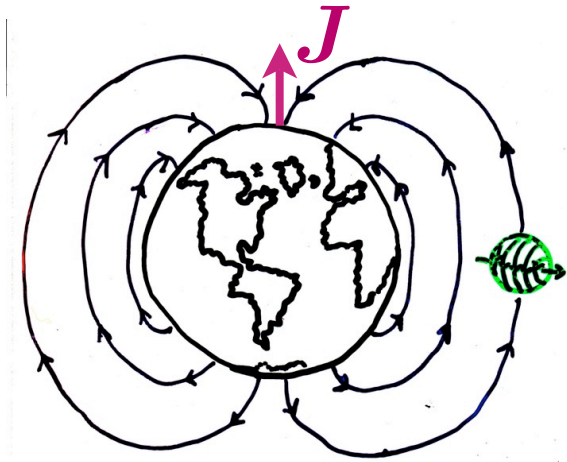


- Head sees feet dragged counter-clockwise
- Feet see head dragged counter-clockwise

negative-vorticity
vortex lines

Frame Dragging When Gravity is Weak

- **Frame-Dragging Vortex Line:** Streamline of eigenvector n of B_{jk}
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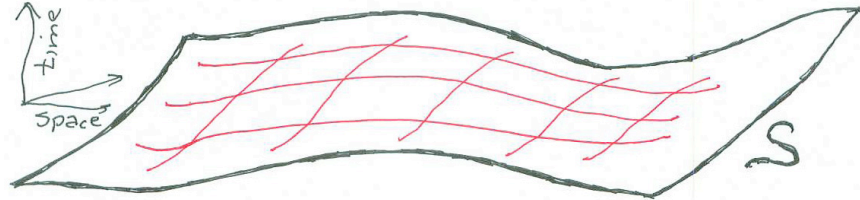


- Head sees feet dragged clockwise
- Feet see head dragged clockwise

positive-helicity
vortex lines

Frame Dragging and Tidal Pull When Gravity is Strong

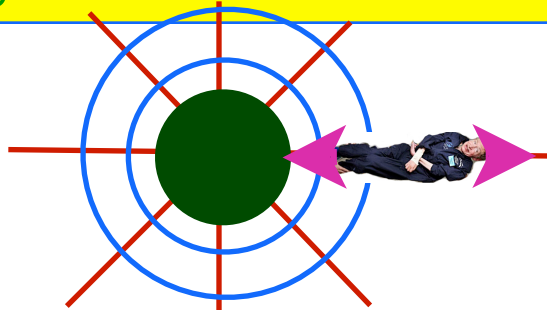
- Slice spacetime into space plus time



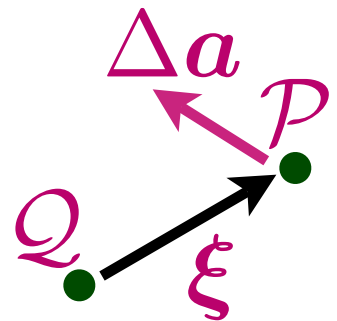
- Observers who move orthogonal to space slices S decompose the spacetime curvature (vacuum Riemann tensor) into “electric” part \mathcal{E}_{jk} and “magnetic” part \mathcal{B}_{jk}
- \mathcal{B}_{jk} describes differential frame dragging - with *vortex lines* and their *vorticities*

- \mathcal{E}_{jk} describes tidal accelerations $\Delta a_j = -\mathcal{E}_{jk} \xi^k$

- \mathcal{E}_{jk} : *tendex* lines and their *tendicity*

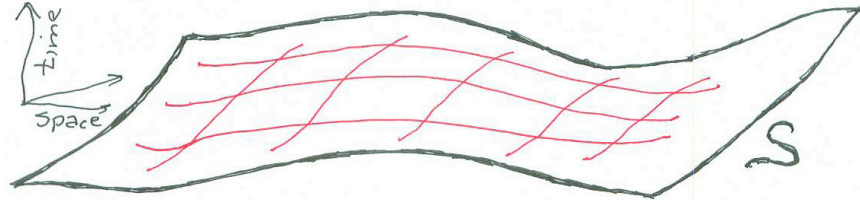


$\mathcal{E}_{nn} < 0$ tidal stretch



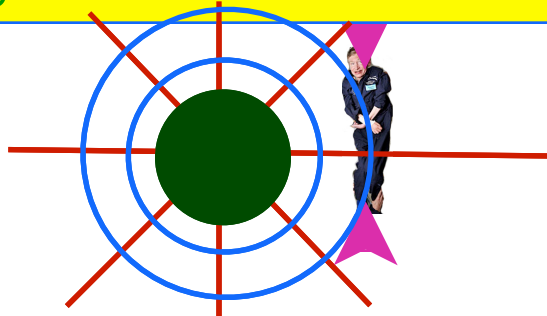
Frame Dragging and Tidal Pull When Gravity is Strong

- Slice spacetime into space plus time

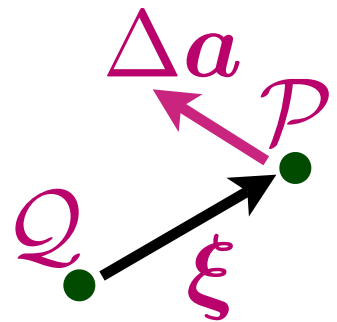


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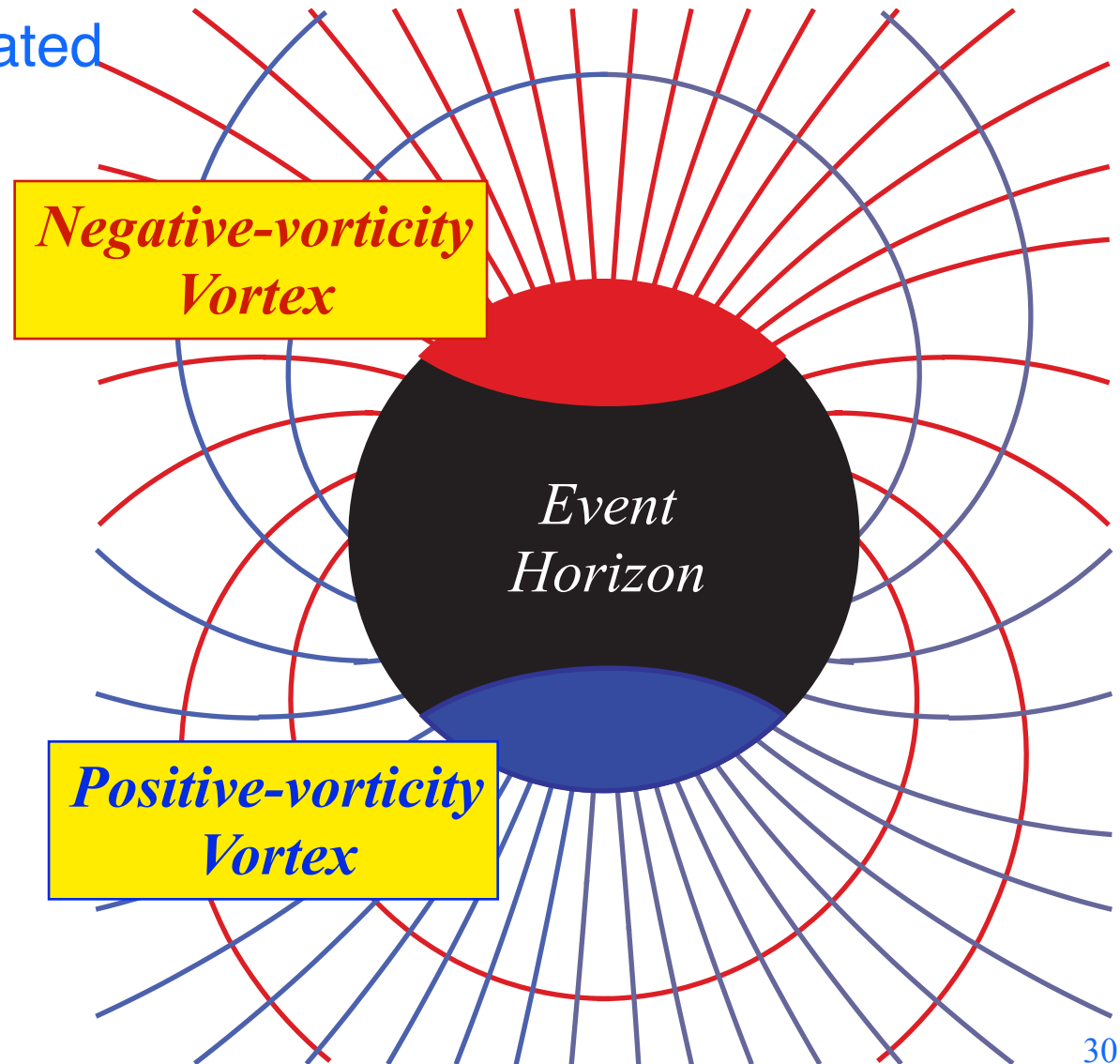


$\mathcal{E}_{nn} > 0$ tidal squeeze



Quiescent, Spinning (Kerr) Black Hole

- *Frame-Dragging Vortex:*
Region of concentrated vorticity



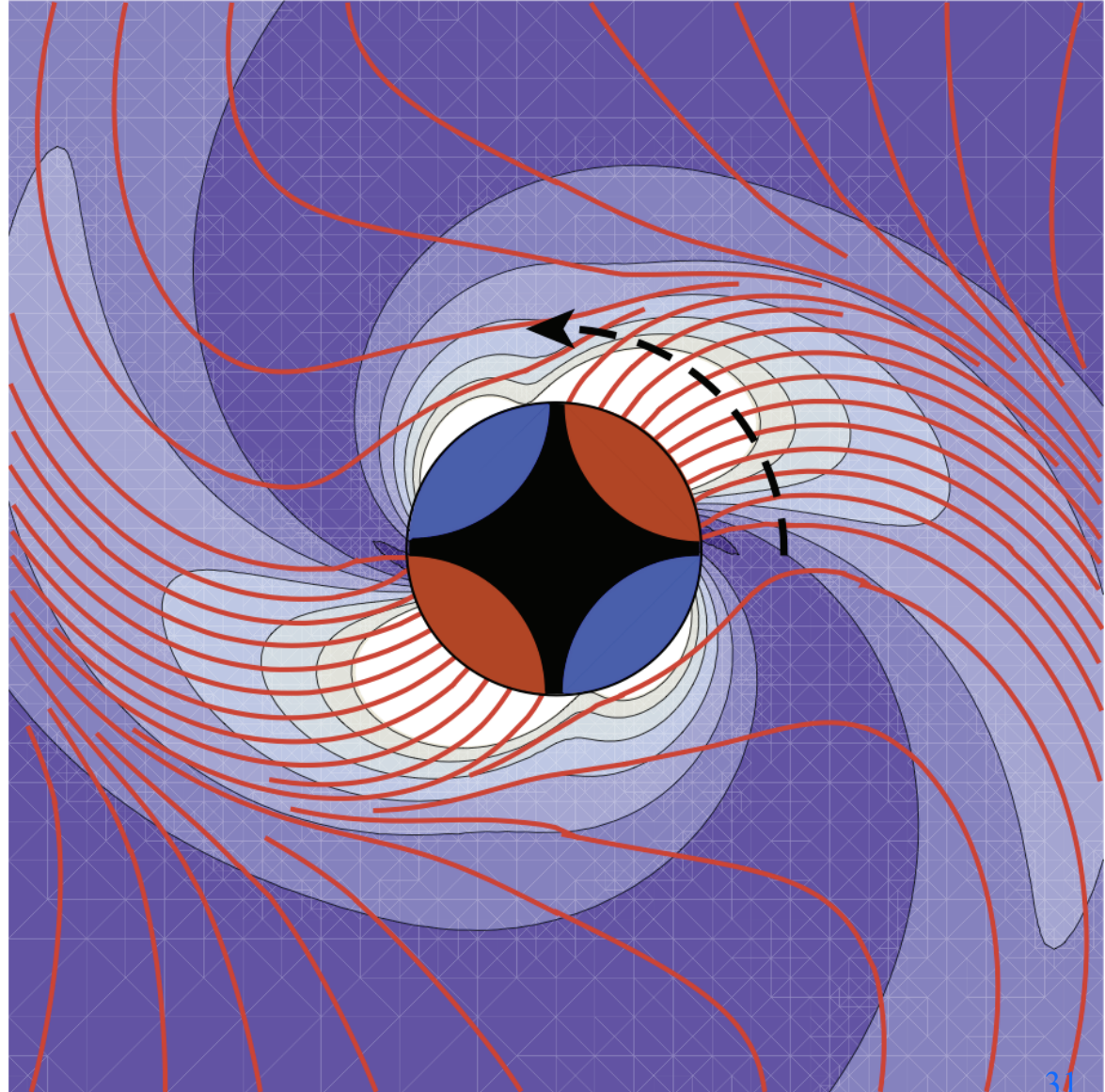
Normal-mode Vortexes Traveling Around a Nonrotating hole

$l=m=2$, *odd-parity*
normal mode

$$\omega = (0.747 - i 0.178) / 2M$$

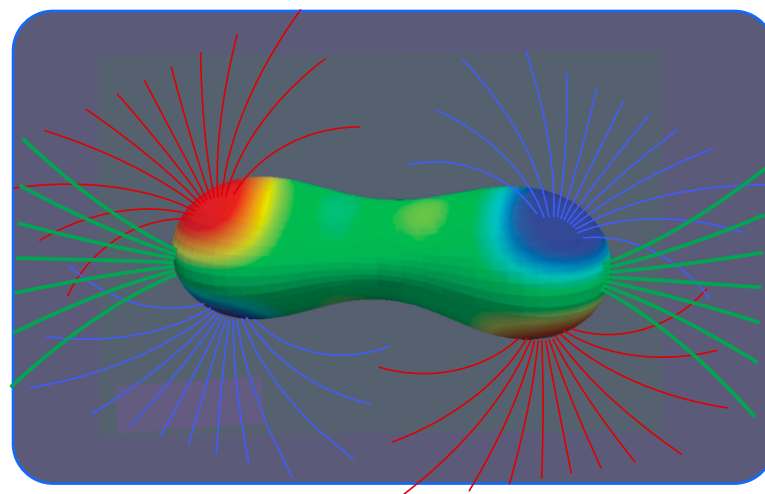
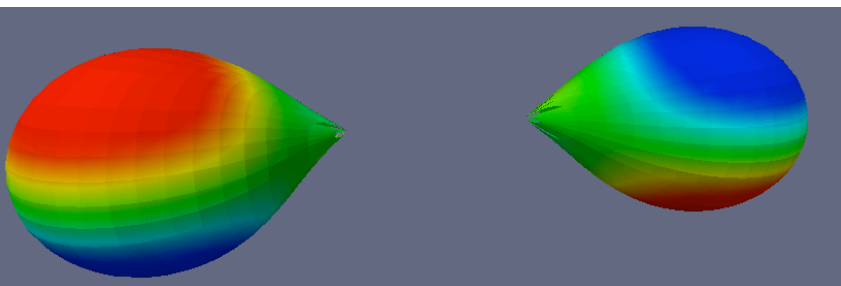
[Chandrasekhar & Detweiler]

- Vortex transitions smoothly into gravitational wave trough
- Slow slippage of vortex lines relative to vortex



Mergers of Spinning Holes

Simulations by Cornell/Caltech/CITA Group



Tendex
lines

Vortex lines

- *For Relativists: On event horizon:*

$$\Psi_2 = \mathcal{E}_{NN} + i\mathcal{B}_{NN}$$

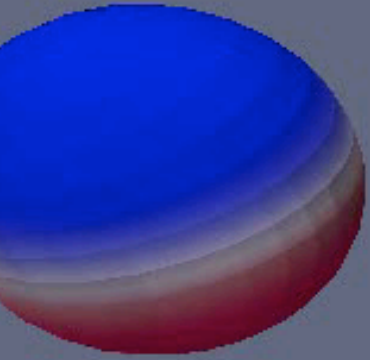
$$\mathcal{K} = \mathcal{R} + i\chi = -\Psi_2 + \mu\rho - \lambda\sigma$$

Complex
curvature

numerically small

Head-On Collision with Transverse Spins

- Keith Matthews, Geoffrey Lovelace, Mark Scheel



Vortexes



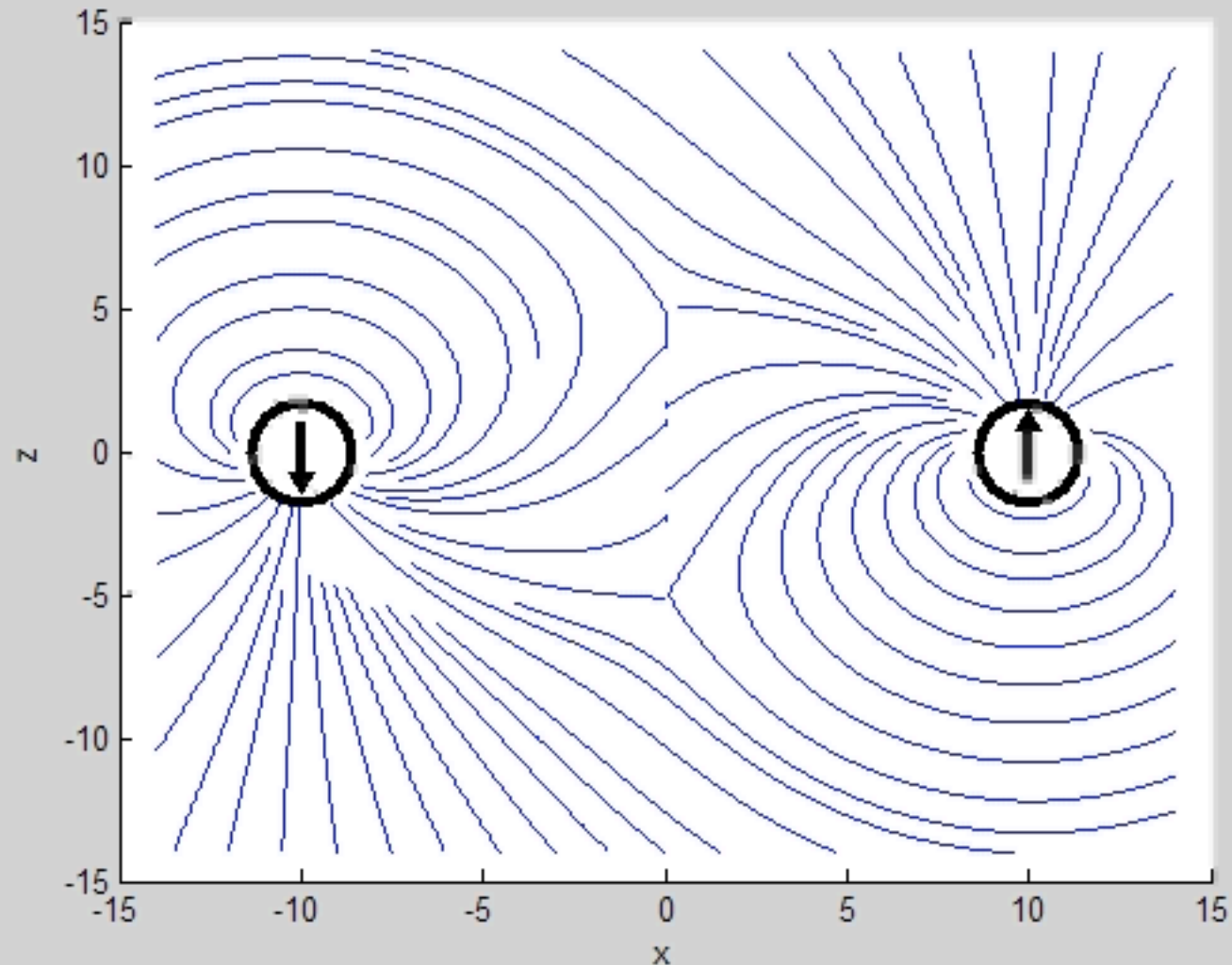
Tendexes



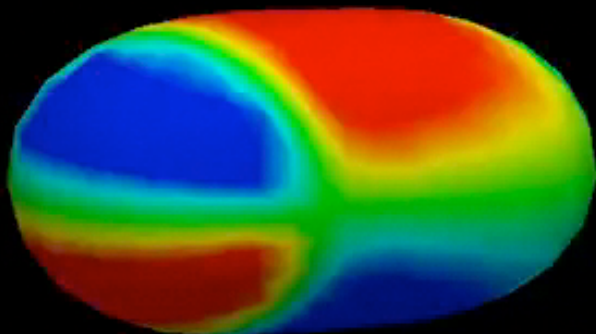
Time: 50.0


Head-On Collision with Transverse Spins

- Evolution of Negative Vorticity Vortex Lines During Inspiral
 - » David Nichols

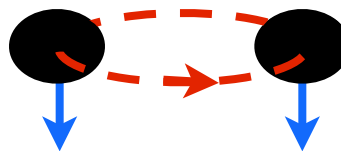


Anti-Aligned Inspiral and Merger




Time: 3483.0

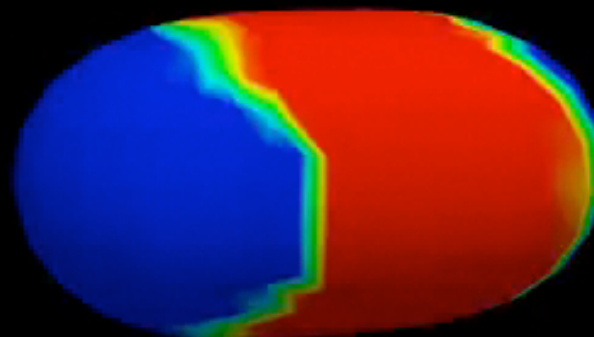
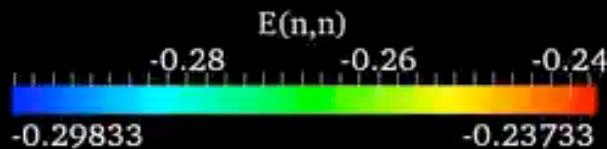
Vortexes



Spin 0.95

- Lovelace, Scheel, Szilagyi

**Common Apparent Horizon
just after merger**

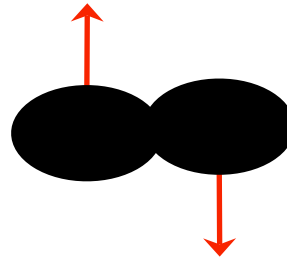
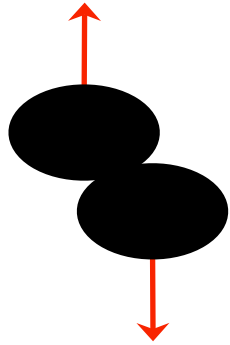
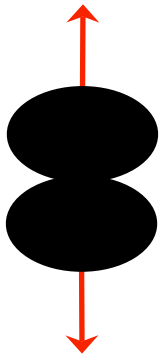


Tendexes

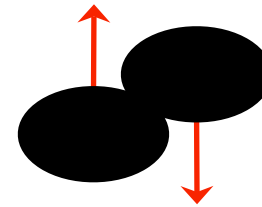
Extreme-Kick Configuration



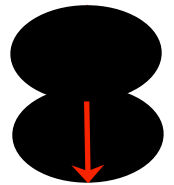
Looking Down from Above



WANT

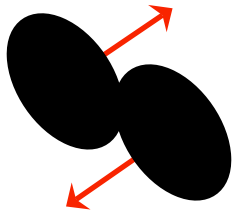


HAVE

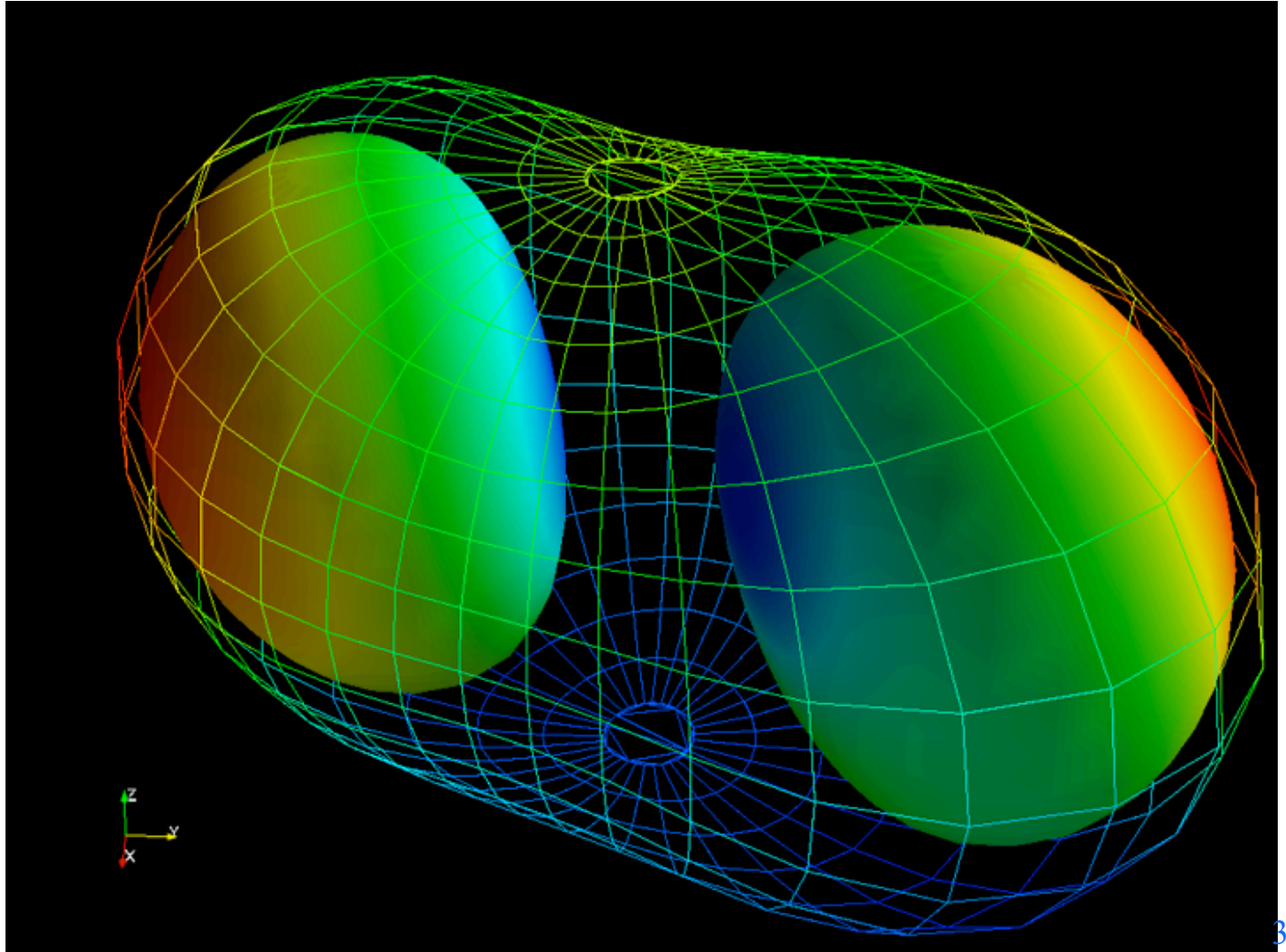


“Common” Apparent Horizon Surrounding Initial Holes’ MOTS

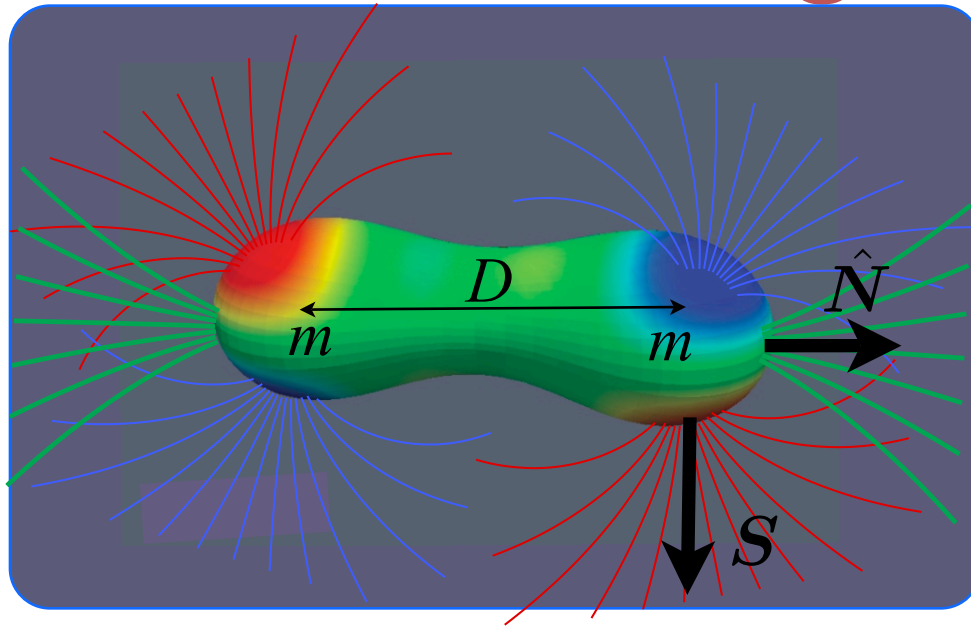
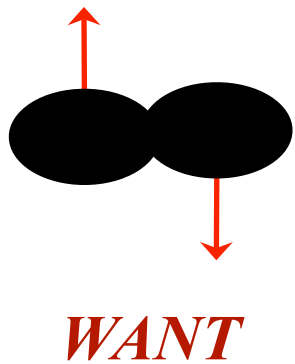
- Rob Owen [Cornell]



HAVE



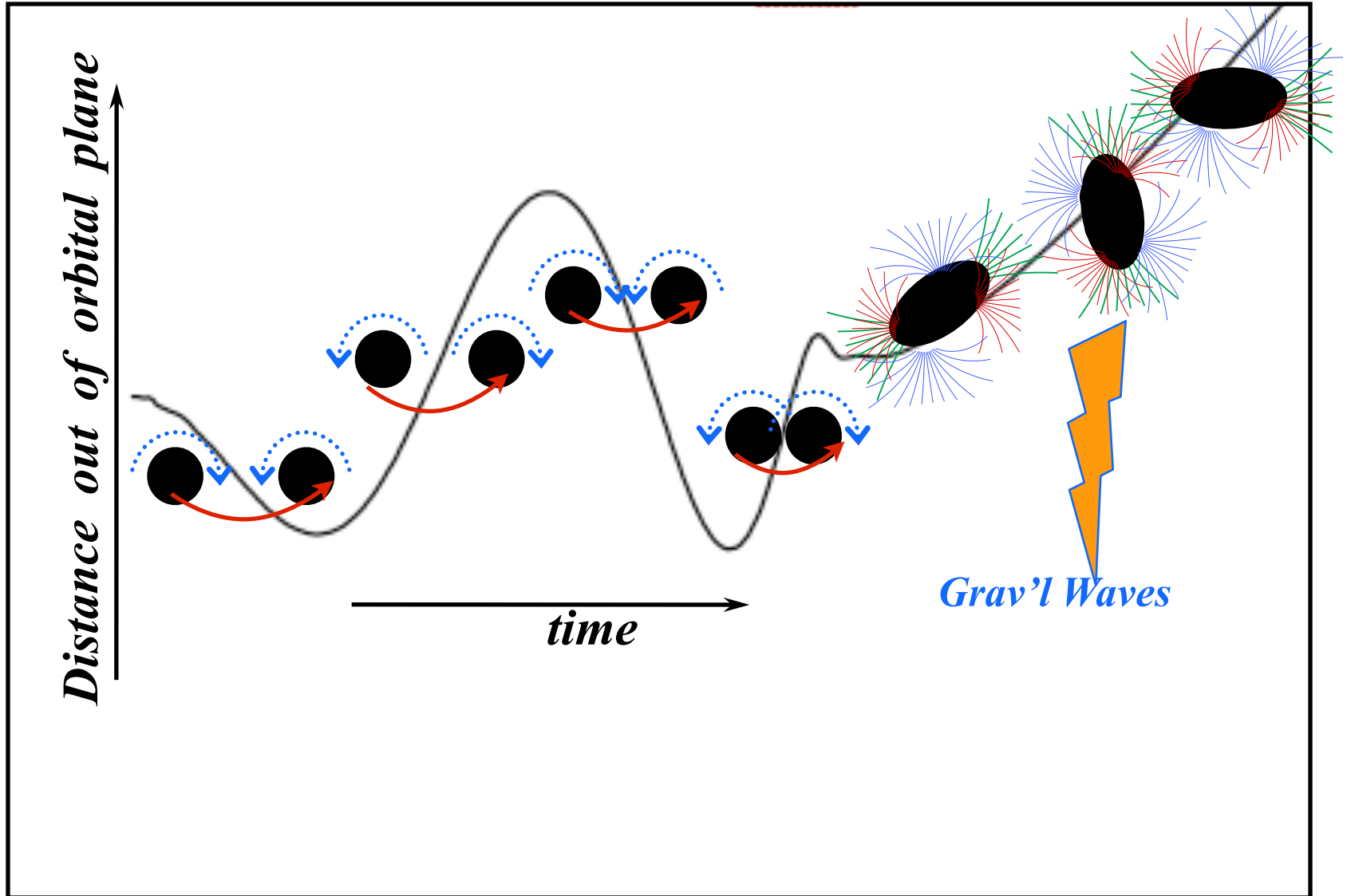
Extreme-Kick Configuration



- Vortexes (S) rotate at Ω_V , Tendexes (N) at Ω_T
- Mass quadrupole GW's beat against current quadrupole GWs

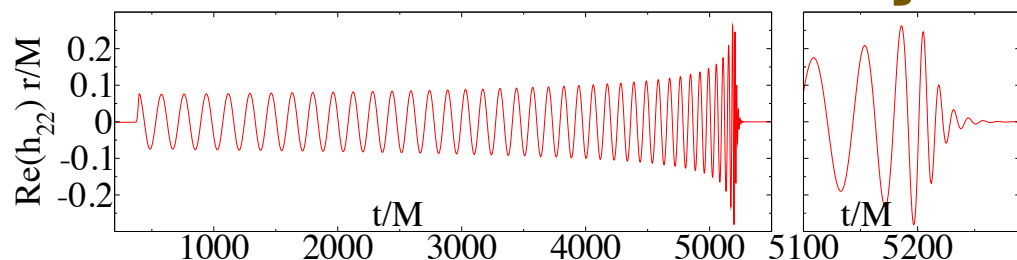
$$\begin{aligned} \frac{dP}{dt} &= -\frac{16}{45} \ddot{\mathbf{I}} \times \ddot{\mathbf{S}} \\ &= -\frac{64}{45} m^2 D^3 S \Omega_T^3 (\Omega_T + \Omega_V)^3 \hat{N} \times S \end{aligned}$$

Extreme Kick



Unreasonable Simplicity of Waveforms

Unreasonable Effectiveness of PN Theory



- **Waveform simplicity:**

- » gravitational waves are produced by dynamics of tendexes and vortexes
- » that dynamics appears to be remarkably simple -
- » perhaps because the near zone is so thin (from $r \sim 2M$ to $r \sim 4M$)

- **Effectiveness of PN Theory (and Close limit Approximation)**
[Will's Talk]

- » perhaps because these approximations capture well the dynamics of the tendexes and vortexes

Conclusion

- Black Holes show an amazing richness of physics
- Numerical Relativity has become a powerful tool
- Gravitational Waves will bring this rich physics into the realm of observations
- A new golden age of black hole research